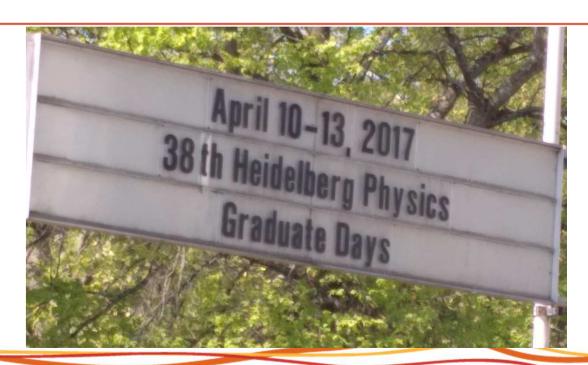




# PHYSICS WITH A FEW ATOMS, PHOTONS AND PHONONS





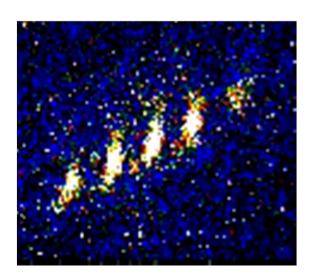


# PHYSICS WITH A FEW ATOMS, PHOTONS AND PHONONS

Take home msg

Problems to solve

Assumptions / derivation steps







# PHYSICS WITH A FEW ATOMS, PHOTONS AND PHONONS

- 1. Ion trap Penning and Paul trap
- 2. Ion light interaction
- 3. Experiments with Penning traps
- 4. Experiments with Paul traps





## **ION TRAP**

#### Content

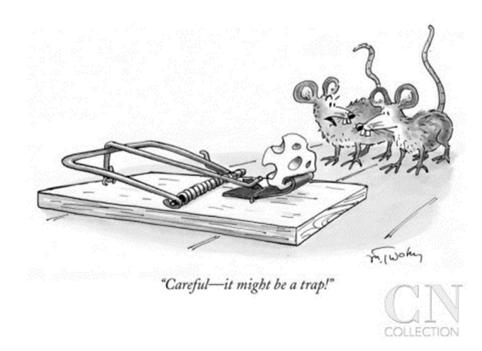
- Usefulness
- 2. Definition and principle of operation
- 3. Types of ion traps
- 4. Stability and equation of motion
- 5. Trap dynamics at low temperature





#### **Definition**

Anything is a trap if you're not careful







#### **Application**

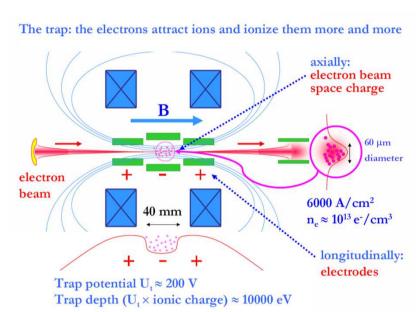
**Particle trap** 

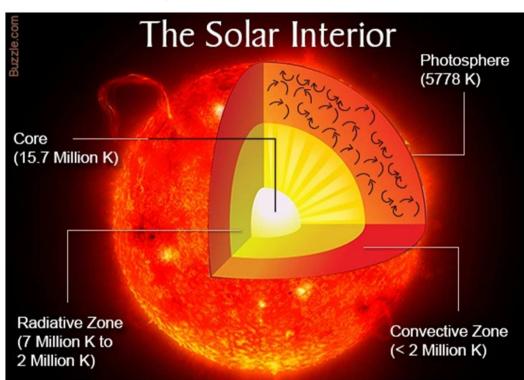
EBIT, super-EBIT etc., Lab. Astrophysics Paul trap Plasma Physics Penning trap Optical lattices, BEC, Condensed matter Paul trap FTICR Penning trap **Biochemistry** Ultra cold chemistry Optical lattices, Paul traps Optical lattices, Paul traps Metrology Paul trap: linear QI & QC and planar Quantum optics Paul trap, dipole trap etc. **Nuclear Physics** Paul, Penning, Optical Particle Physics Paul trap





## Application (Astro. and plasma)





Problem 1a: What is the electron binding energy of hydrogen-like Uranium? Z=92

Priv.-Doz. Dr. José Ramon Crespo López-Urrutia, MPIK





# **Application**

**Particle trap** 

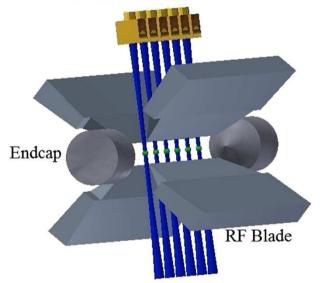
Lab. Astrophysics	EBIT, super-EBIT etc., Paul trap
— Plasma Physics	Penning trap
— Condensed matter	Optical lattices, BEC, Paul trap
Biochemistry	FTICR Penning trap
Ultra cold chemistry	Optical lattices, Paul traps
Metrology	Optical lattices, Paul traps
QI & QC	Paul trap: linear and planar
Quantum optics	Paul trap, dipole trap etc.
Nuclear Physics	Paul, Penning, Optical
Particle Physics	Paul trap

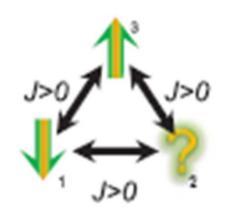


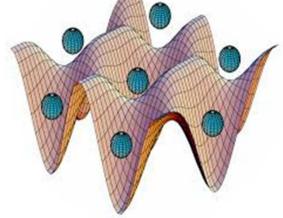


## Application (condensed matter)

Piezoelectric micromirror actuator







Frustration in spin system

K. Kim et al. NJP 13, 105003 (2011)

Super fluidity and conductivity

T. Dutta et al. PRL 111, 170406 (2013)





## **Application**

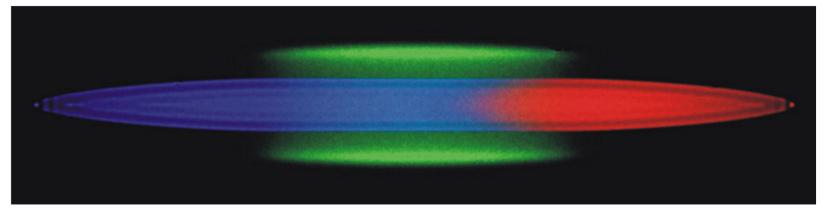
**Particle trap** 

Lab. Astrophysics	EBIT, super-EBIT etc., Paul trap
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Particle Physics	Paul trap





# Application (ultra cold chemistry)



Molecular ions at mK temperature

How was life formed?

Problem 1b: Is this temperature good enough for cold chemistry?

Phys. Rev. Lett. 97, 243005 (2006)





# **Application**

**Particle trap** 

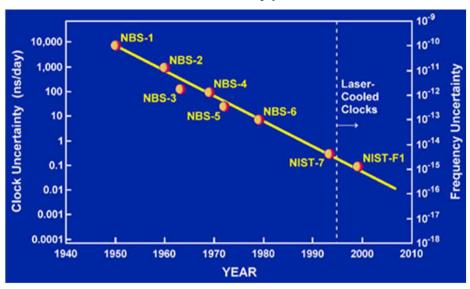
Lab. Astrophysics	EBIT, super-EBIT etc., Paul trap
— Plasma Physics	Penning trap
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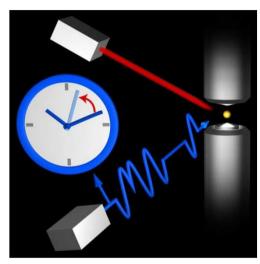


## Application (metrology)

9,192,631,770 Hertz hyperfine transition 133Cs



$$\frac{\delta v}{v} = 10^{-14}$$



N. Huntemann et al. Phys. Rev. Lett. **116**, 063001 (2016)

$$\frac{\delta \nu}{\nu} = 10^{-18}$$

Problem 1c: with these accuracy how fast or slow will your clock run in day?





# **Application**

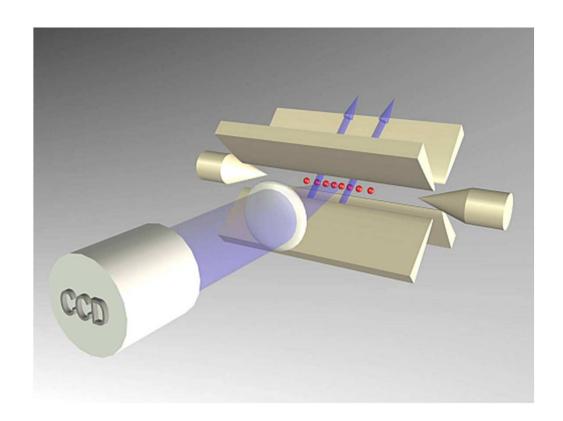
**Particle trap** 

Lab. Astrophysics	EBIT, super-EBIT etc., Paul trap
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Nuclear Physics	Paul, Penning, Optical
Particle Physics	Paul trap





## Application (QC, QO, QI)



IQOQI, University Innsbruck





#### **Application**

**Particle trap** 

EBIT, super-EBIT etc., Lab. Astrophysics Paul trap Plasma Physics Penning trap Optical lattices, BEC, Condensed matter Paul trap FTICR Penning trap **Biochemistry** Ultra cold chemistry Optical lattices, Paul traps Optical lattices, Paul traps Metrology Paul trap: linear QI & QC and planar Quantum optics Paul trap, dipole trap etc. **Nuclear Physics** Paul, Penning, Optical Particle Physics Paul trap





### Application (constituents of the Universe)

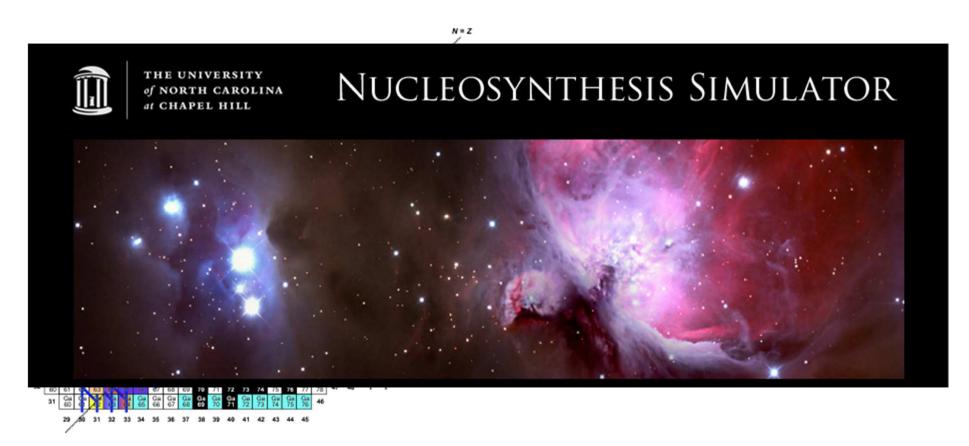








## Application (Nuclear physics)



Blaum, Klaus et al. Phys. Scripta T152 014017 (2013)





# **Application**

**Particle trap** 

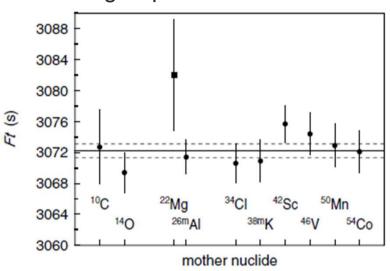
Lab. Astrophysics	EBIT, super-EBIT etc., Paul trap
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Nuclear Physics	Paul, Penning, Optical
Particle Physics	Paul trap





## Application (Particle physics)

Penning trap mass measurements

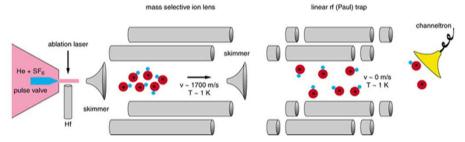


Conserved vector current hypothesis CKM unitarity

M. Mukherjee et al. PRL 93, 150801(2004)

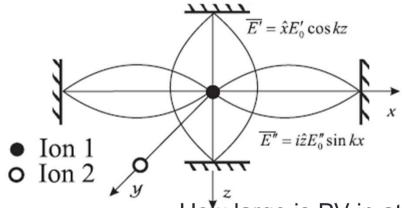
Is the SM correct?

Paul trap laser spectroscopy



Does electron has a permanent EDM?

EDM ion trap experiment, JILA NIST (E. Cornell)



How large is PV in atoms? CQT/NUS, UW, KVI etc.

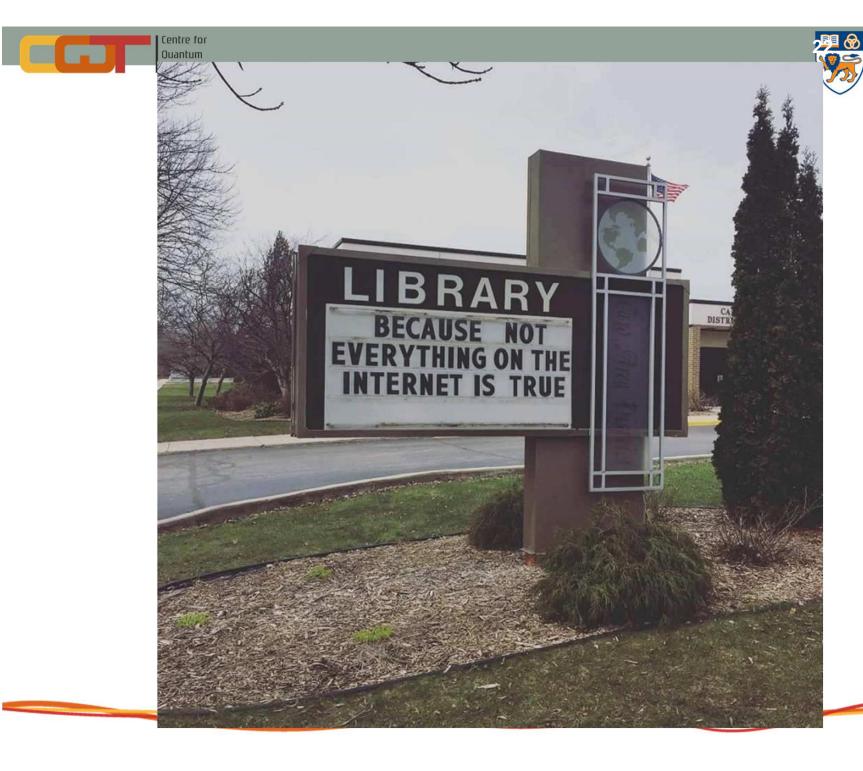




#### **Definition**

Confinement of particle/s in a small volume on space





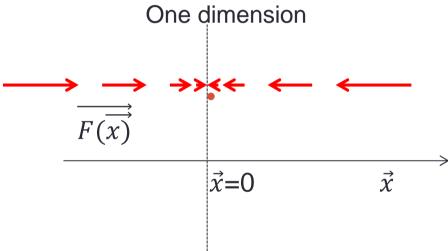
National University of Singapore





#### Principle of operation (force argument)

 To confine a particle in one dimension the required force should be proportional to the displacement but opposite in direction



$$\vec{F}(x) = -\alpha \vec{x}$$

$$-\frac{dV}{dx} = -\alpha x$$

$$V = \frac{\alpha}{2}x^2$$





## Trapping charged particle

Continuing to 2D and applying same logic we get

$$V(x,y) = \alpha x^2 + \beta y^2$$

However there is no charge so, *Laplace* condition needs to be fulfilled  $\nabla^2 V = 0$ 

So,

$$V(x,y) = \alpha x^2 - \beta y^2$$

Hence

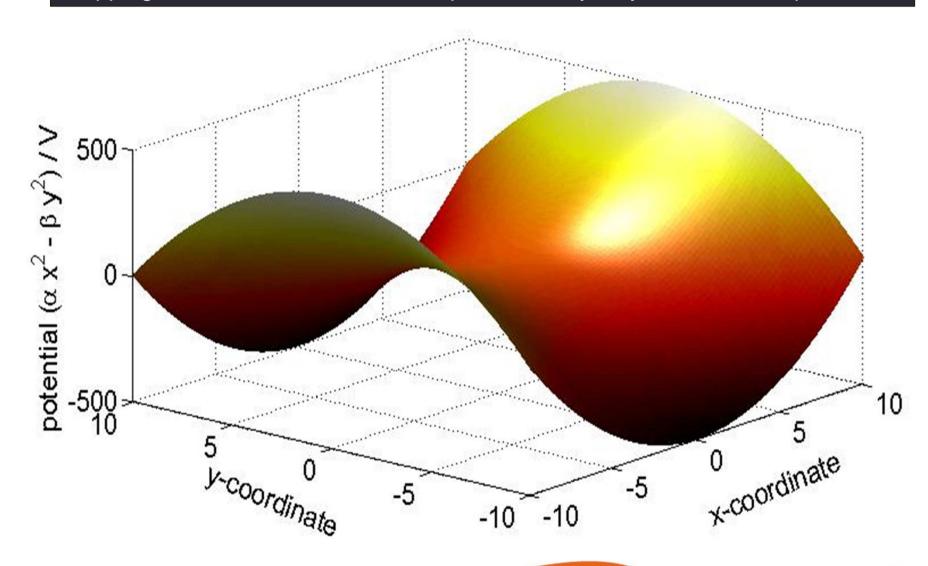
$$\vec{F}(x,y) = -\alpha x \hat{x} + \beta y \hat{y}$$

Anti-trapping





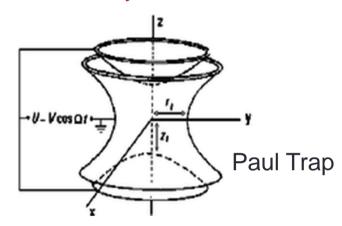
#### Trapping in more than 1D cannot be performed by only static electric potentials



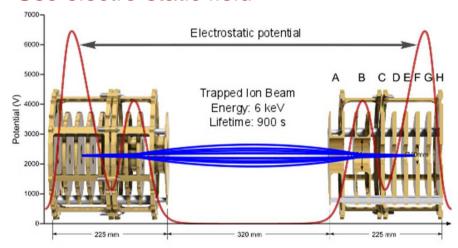




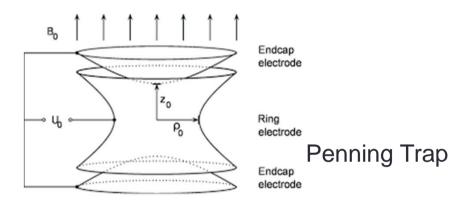
#### Use electro-dynamic field



#### Use electro-static field



#### Use electro-static + magneto-static field



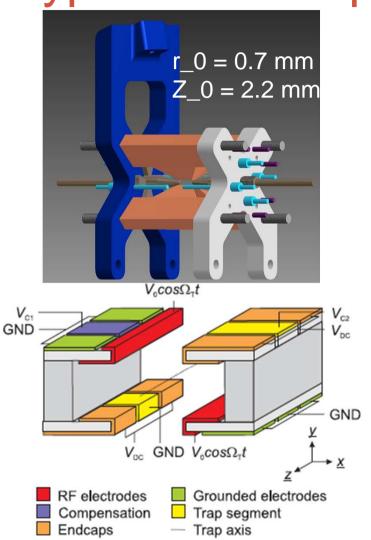
#### Sources:

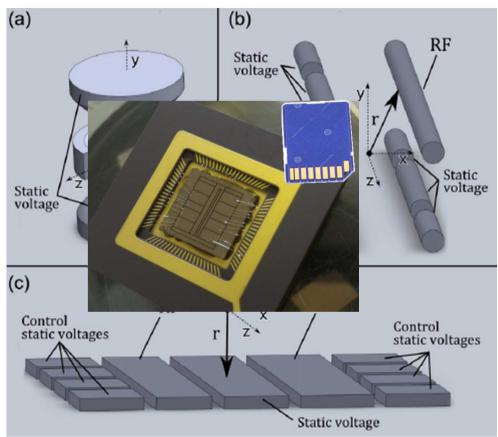
Mass Spectrom. Lett. 2015 Vol. 6, No. 4, 112–115 (2015) <a href="https://www.mpi-hd.mpg.de/blaum/high-precision-ms/ptms/basics.en.html">https://www.mpi-hd.mpg.de/blaum/high-precision-ms/ptms/basics.en.html</a>

Electro-static Trap









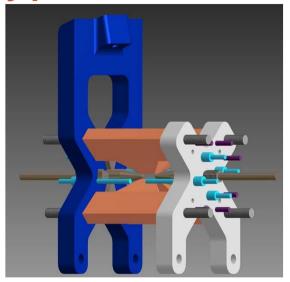
https://arxiv.org/abs/1106.5013

NJP, 8, 232 (2006)

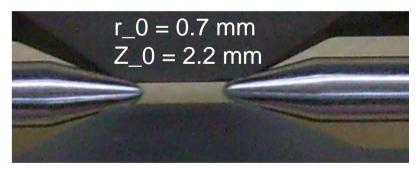
http://coldiongroup.wixsite.com/index

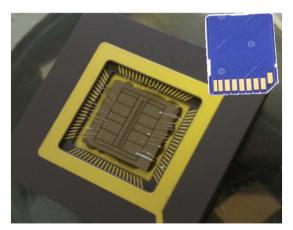












http://coldiongroup.wixsite.com/index

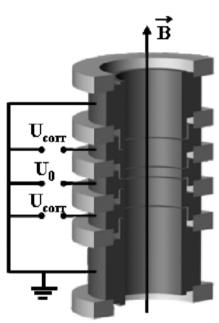












Physics world (2013) MPIK Heidelberg ISOLTRAP, CERN Can. J. Physics (2010)





#### Principles of trap operation ~ History

"We never experiment with just one electron or atom or small molecule. In thought experiments we sometime assume that we do, this invariably entails ridiculous consequences ....In the first place it is fair to state that we are not experimenting with single particle anymore than we can raise Ichthyosauria in the zoo" ------ 1952

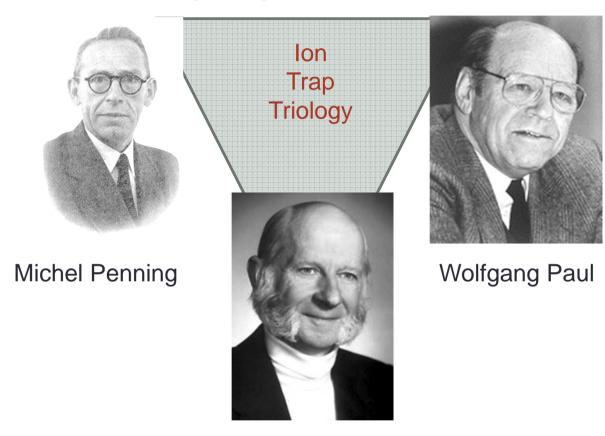
"Professor Dr **Wolfgang Paul**, University of Bonn, Federal Republic of Germany and Professor **Hans G. Dehmelt**, University of Washington, USA, have introduced and developed the *ion trap technique* which has made it possible to study a single electron or a single ion with extreme precision." ----- 1989

"Serge Haroche and David J. Wineland have independently invented and eveloped methods for measuring and manipulating individual particles while preserving their quantum-mechanical nature, in ways that were previously thought unattainable." ------2012





## Principles of trap operation ~ History



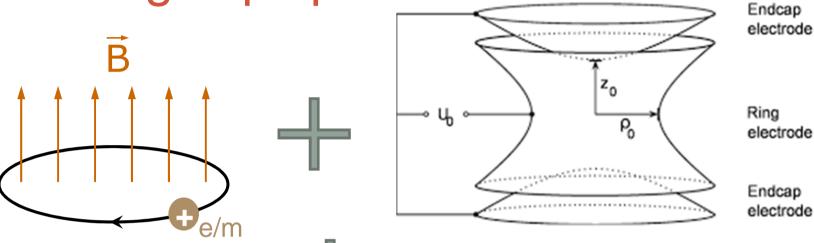
Hans G. Dehmelt

Michael H. Holzscheiter, Physica Scripta. Vol. T59, 69-76, 1995





#### Penning trap operation



Static magnetic field



Static electric field

Lorentz force:

$$\vec{F} = e(\vec{E} + \vec{v} \times \vec{B})$$

Transform to another inertial frame

$$\vec{v} = \vec{v'} + \vec{E} \times \frac{\vec{B}}{B^2}$$

$ec{E}$	electric field
$\overrightarrow{B}$	magnetic field
$ec{v}$	velocity of charge e





#### Penning trap operation ~ principle

$$m\frac{d\vec{v}}{dt} = e(\vec{E} + \vec{v} \times \vec{B})$$

On substitution

$$m\frac{d\overrightarrow{v'}}{dt} = e(\overrightarrow{E} + \left(\overrightarrow{v'} + \overrightarrow{E} \times \frac{\overrightarrow{B}}{B^2}\right) \times \overrightarrow{B}$$

$$\Rightarrow m \frac{d\overrightarrow{v'}}{dt} = e\left(\overrightarrow{v'} \times \overrightarrow{B}\right)$$

$$(\vec{A} \times \vec{B}) \times \vec{C}$$

$$= -\vec{C} \times (\vec{A} \times \vec{B})$$

$$= (\vec{C} \cdot \vec{A}) \vec{B}$$

$$- (\vec{C} \cdot \vec{B}) \vec{A}$$

$$\vec{E} \cdot \vec{B} = 0$$

With this newly defined velocity, the eq. of motion for a charged particle without electric field is recovered. This means that the role of electric field can be substituted by a  $(\vec{E} \times \vec{B})$  field.

Therefore balance of force leads to:





#### Penning trap operation ~ principle

#### Pure cyclotron motion

Motion of a charged particle in pure magnetic field

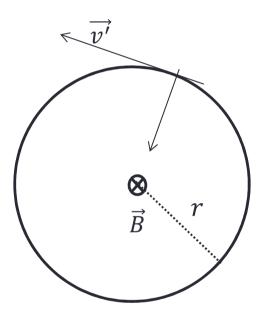
$$\frac{mv'^2}{r} = ev'E$$

$$\Rightarrow \omega' = \frac{eB}{m}$$

Magnetic field strength



$$\omega' = \frac{eB}{m}$$
Mass



Problem 1: Calculate the cyclotron frequencies of an electron and a singly charged calcium atom in an uniform magnetic field of 1T?

Problem 2: Suppose both electron and positron are confined in a same uniform magnetic field. What would be the difference in their motion?





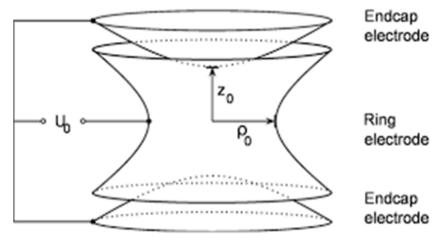
#### Penning trap operation ~ equation of motion

The electrostatic part is obtained from the voltages applied to the electrodes which are hyperbolic in shape

$$V(\rho, z) = \frac{\frac{U_0}{2} \left(z^2 - \frac{\rho^2}{2}\right)}{d_0^2}$$

With characteristic dimension

$$d_0^2 = \frac{1}{2} \left( z_0^2 + \frac{\rho_0^2}{2} \right)$$



This type of potential originates by following the surface of the electrodes which forms two hyperbolas namely,

$$z^2 = \left(z_0^2 + \frac{\rho^2}{2}\right)$$

$$z^2 = \frac{1}{2}(\rho_0^2 - \rho^2)$$

The ring electrode

The endcap electrode





#### Penning trap operation ~ equation of motion

Now combine electro-static and magneto-static parts:

$$m\frac{d\vec{v}}{dt} = e(\vec{E} + \vec{v} \times \vec{B})$$

since

$$\dot{z}\hat{z} \times \vec{B} = 0$$

$$\vec{r}(x, y, z) = x\hat{x} + y \hat{y} + z \hat{z}$$
$$= \rho \hat{\rho} + z \hat{z}$$

$$ec{\dot{r}} = ec{v} = \dot{
ho}ec{
ho} + \dot{z}\hat{z}$$

The <u>z-motion</u> or axial motion is decoupled/independent from radial/transverse motion denoted by  $\rho$ .

$$\ddot{z} + \omega_Z^2 z = 0 \quad \dots (1)$$

$$\omega_z^2 = \frac{eU_0}{md_0^2}$$

The  $\rho$ -motion or radial motion is governed by:

$$\ddot{\rho} = \frac{e}{m} \left( \vec{E}_{\vec{\rho}} + \dot{\vec{\rho}} \times \vec{B} \right)$$





#### Penning trap operation ~ equation of motion

$$\ddot{\rho} = \frac{e}{m} \left( \vec{E}_{\vec{\rho}} + \dot{\vec{\rho}} \times \vec{B} \right)$$

The radial component of the electric field can be obtained as:

$$\overrightarrow{E_{\vec{\rho}}} = -\frac{\partial V}{\partial \rho} \hat{\rho} = \frac{U_0 \vec{\rho}}{2d_0^2} = \frac{1}{2} \frac{m}{e} \ \omega_z^2 \ \vec{\rho}$$

The last step is substitution

$$\omega_z^2 = \frac{eU_0}{md_0^2}$$

The equation of motion can now be re-written as:

$$\ddot{\vec{\rho}} - \omega_c \, \dot{\vec{\rho}} \times \hat{B} \, -\frac{1}{2} \omega_z^2 \, \vec{\rho} = 0 \quad \dots (2)$$

The eigen frequencies obtained by solving eq. (2)

$$\omega_{\pm} = \frac{1}{2} \left( \omega_c \pm \sqrt{\omega_c^2 - 2\omega_z^2} \right)$$

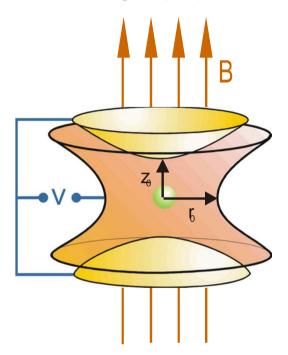
Modified cyclotron frequency :  $\omega_+$ 

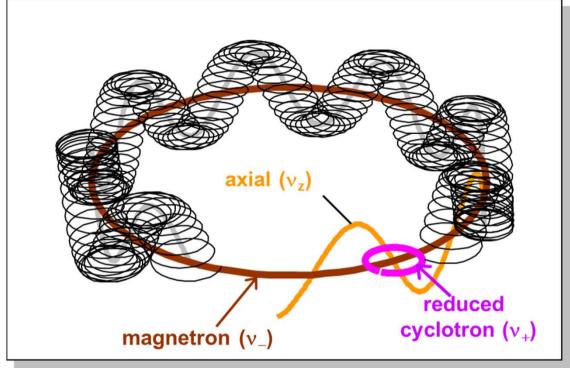
Magnetron frequency :  $\omega_{-}$ 





#### Penning trap operation ~ equation of motion





Courtesy: Prof. Dr. K. Blaum, MPIK Heidelberg

### Changing to frequency units

$$\omega = 2\pi \times \nu$$

$$\nu_{\pm} = \frac{\nu_c}{2} \pm \sqrt{\frac{{\nu_c}^2}{4} - \frac{{\nu_z}^2}{2}}$$

Invariance Theorem

$$\nu_c = \nu_+ + \nu_-$$





Problem 3: Calculate the modified cyclotron and magnetron frequency of a trapped calcium singly charge ion for 1T magnetic field and a trap of dimensions  $r_0=1\,$  mm and  $z_0=2\,$  mm while endcap voltage of 100V applied to it.?





Re-write the radial velocity components in terms of new radial velocities as

$$\overrightarrow{V^{\pm}} = \overrightarrow{\dot{\rho}} - \omega_{\mp} \, \hat{z} \times \overrightarrow{\rho}$$

Taking temporal derivative and substituting  $\vec{\ddot{p}}$ 

$$\overrightarrow{\dot{V}^{\pm}} = \omega_{+} \hat{z} \times \overrightarrow{V^{\pm}}$$
 2 de-coupled motion in the radial plane

Now, consider a gauge for the vector potential as  $\vec{A} = \frac{1}{2}\vec{B} \times \vec{\rho}$ , the canonical momentum becomes

$$\vec{p} = m\vec{\dot{\rho}} + (e)\vec{A}$$

$$\vec{p} = \frac{m}{2} \left( \overrightarrow{V^+} + \overrightarrow{V^-} \right)$$

$$\omega_+ + \omega_- = \omega_c = \frac{eB}{m}$$





The commutation relations between the components of  $V^{\pm}$  are to be derived. So we start with  $\rho_k s$  and  $p_l s$ 

$$[\rho_k, p_l] = i\hbar \delta_{kl} \Rightarrow [\rho_k, \dot{\rho_l}] = \frac{i\hbar}{m} \delta_{kl}$$
  $k, l$  are x, y components of  $\rho$ 

$$\left[\dot{\rho_{x}},\dot{\rho_{y}}\right] = \frac{ie\hbar}{m^{2}c} \left[ \frac{\partial A_{y}}{\partial x} - \frac{\partial A_{x}}{\partial y} \right] = -\frac{i\hbar \left| e\vec{B} \right|}{m^{2}c} = -i\frac{\hbar\omega_{c}}{m}$$

Therefore we obtain

$$[V_k^+, V_l^-] = 0$$
 and  $\mp \frac{m}{\omega_+ - \omega_-} [V_x^{\pm}, V_y^{\pm}] = i\hbar$  .....(4)

The radial Hamiltonian

$$H_{\rho} = \frac{1}{2} m \left( \dot{\rho}^2 - \frac{1}{2} \omega_z^2 \rho^2 \right) \Rightarrow \frac{1}{2} m \frac{\omega_+ V^{+2} - \omega_- V^{-2}}{\omega_+ - \omega_-}$$

Magnetron has -ve energy contribution and hence unstable





$$[V_k^+, V_l^-] = 0$$
 and  $\mp \frac{m}{\omega_+ - \omega_-} [V_x^{\pm}, V_y^{\pm}] = i\hbar$  .....(4)

This suggests  $V_x^{\pm}$  and  $V_y^{\pm}$  can be treated as conjugate pair of variables (p,q).

With proper normalization one obtains

$$a_{\pm} = \left[\frac{m}{2\hbar(\omega_{+} - \omega_{-})}\right]^{\frac{1}{2}} \left(V_{x}^{\pm} \mp iV_{y}^{\pm}\right)$$

$$a_{\pm}^{\dagger} = \left[\frac{m}{2\hbar(\omega_{+} - \omega_{-})}\right]^{\frac{1}{2}} \left(V_{x}^{\pm} \pm iV_{y}^{\pm}\right)$$

$$\left[a_{\pm}, a_{\pm}^{\dagger}\right] = 1$$

The Hamiltonian is given as:

$$H_{\rho} = \hbar\omega_{+}\left(a_{+}^{\dagger}a_{+} + \frac{1}{2}\right) - \hbar\omega_{-}\left(a_{-}^{\dagger}a_{-} + \frac{1}{2}\right)$$



# National University of Singapore

#### Penning trap operation ~ Quantum treatment

The radial Hamiltonian is given as:

$$H_{\rho} = \hbar\omega_{+}\left(a_{+}^{\dagger}a_{+} + \frac{1}{2}\right) - \hbar\omega_{-}\left(a_{-}^{\dagger}a_{-} + \frac{1}{2}\right)$$

The axial Hamiltonian is

$$H_z = \frac{P_z^2}{2m} + \frac{m\omega_z^2 z^2}{2}$$
 with  $[z, P_z] = i\hbar$ 

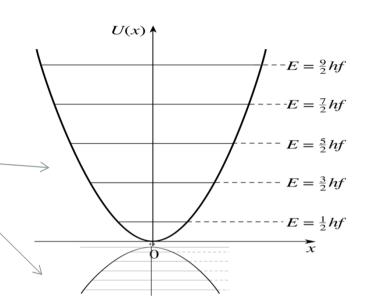


$$a_z^{\dagger} = \left[\frac{m\omega_z}{2\hbar}\right]^{\frac{1}{2}} z - i \left[\frac{1}{2m\hbar\omega_z}\right]^{\frac{1}{2}} P_z$$

$$a_z = \left[\frac{m\omega_z}{2\hbar}\right]^{\frac{1}{2}} z + i \left[\frac{1}{2m\hbar\omega_z}\right]^{\frac{1}{2}} P_z$$

The axial Hamiltonian is:

$$H_z = \hbar \omega_z \left( a_z^{\dagger} a_z + \frac{1}{2} \right)$$







The Hamiltonian is given as:

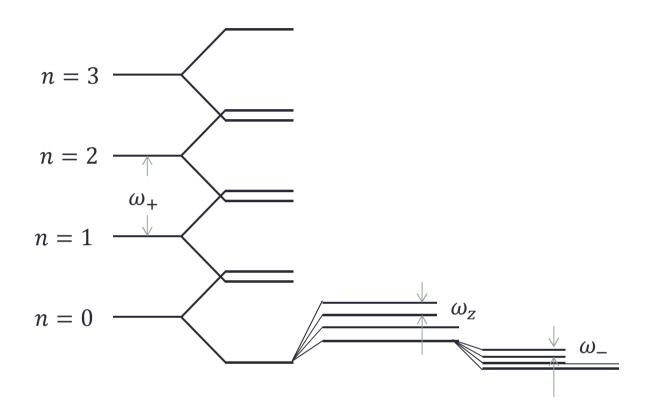
$$H = \hbar\omega_{+}\left(a_{+}^{\dagger}a_{+} + \frac{1}{2}\right) + \hbar\omega_{z}\left(a_{z}^{\dagger}a_{z} + \frac{1}{2}\right) - \hbar\omega_{-}\left(a_{-}^{\dagger}a_{-} + \frac{1}{2}\right)$$

Problem 4: Compare the modified cyclotron, axial and magnetron frequencies of a trapped proton and an electron for 1T magnetic field and a trap of dimensions  $r_0 = 1$  mm and  $z_0 = 2$  mm while endcap voltage of 100V applied to it?

Problem 5: Compare the energies in the above three oscillators (for proton and electron) to that of electrode thermal energy kept at 4.2K. Do you need to treat the axial motion quantum mechanically?





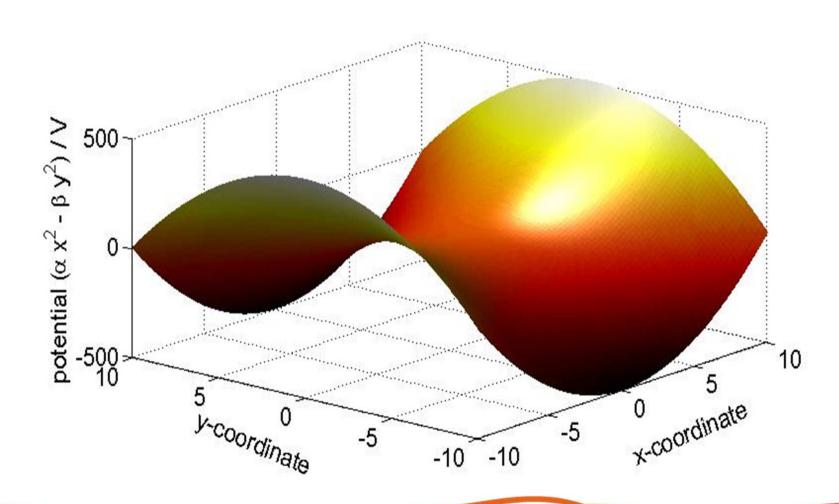


Energy level structure of an single spin-1/2 particle in a Penning trap





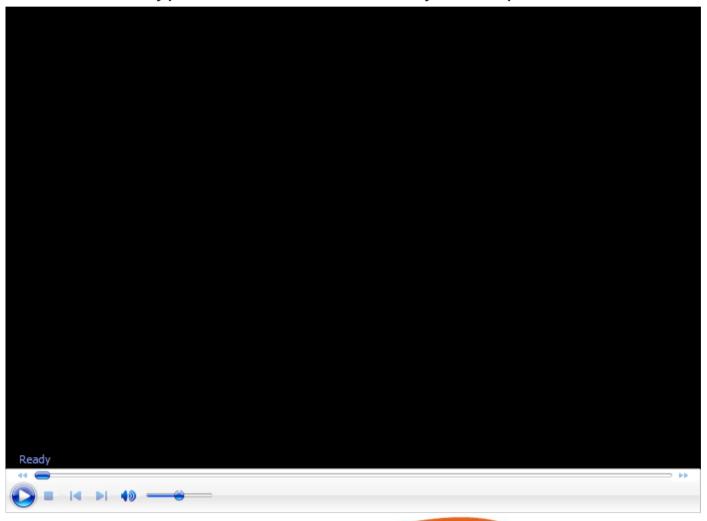
Hyperbolic electrode and static potential







Hyperbolic electrode and dynamic potential







In 3D the potential for confining charged particle needs to be time varying:

$$V(x, y, z) = \frac{U_o}{2}(\alpha x^2 + \beta y^2 + \gamma z^2) + \frac{V_o cos\Omega t}{2}(\alpha' x^2 + \beta' y^2 + \gamma' z^2)$$

In confining region free of source Laplace condition  $\nabla^2 V = 0$ :

$$\alpha + \beta + \gamma = 0$$
 and  $\alpha' + \beta' + \gamma' = 0$ 

This condition can be achieved in different ways leading to differing geometry of Paul traps

Hyperbolic Paul trap

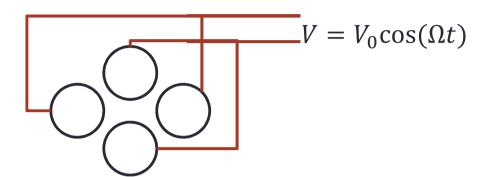
$$\alpha = \beta = \gamma = 0$$
 and  $\alpha' + \beta' = -\gamma'$ 

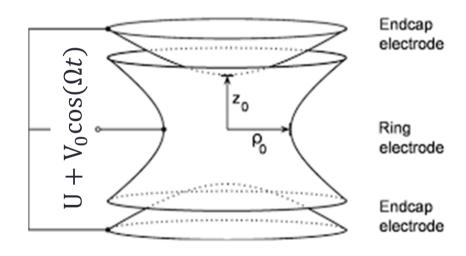
Linear Paul trap

$$-(\alpha + \beta) = \gamma > 0$$
 and  $\alpha' = -\beta'$ ,  $\gamma' = 0$ 

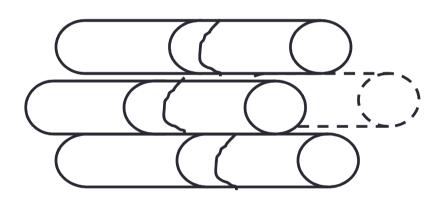












Linear Paul trap





#### Linear traps have the following advantages:

- 1. Ions are easily accessible by lasers
- 2. Can form linear chain of ions
- 3. Individual addressing of ions possible
- 4. Easy to machine





$$\ddot{x} = -\frac{Ze}{m}\frac{\partial V}{\partial x} = -\frac{Ze}{m}\left(U_0\alpha + V_0\alpha'\cos\Omega t\right)x$$

Substitute

$$\zeta = \frac{\Omega t}{2} \ a_x = \frac{4ZeU_0\alpha}{m\Omega^2} \ q_x = -\frac{2ZeV_0\alpha'}{m\Omega^2}$$

$$\ddot{x} + (a_x - 2q_x \cos 2\zeta)x = 0$$

Mathieu equation

It is second order differential equation similar to simple harmonic oscillator equation but with time dependent frequency term





Stable solution for the Mathieu equations

$$x(\zeta) = Ae^{\{i\beta_x\zeta\}} \sum_{\{n=-\infty\}}^{\infty} C_{2n} e^{\{i2n\zeta\}} + Be^{\{-i\beta_x\zeta\}} \sum_{\{n=-\infty\}}^{\infty} C_{2n} e^{\{-i2n\zeta\}}$$

Depends on initial condition

Depends on  $a_x$ ,  $q_x$  not on boundary condition

Using the solution in the equation of motion one obtains the following relation:

$$C_{2n+2} - D_{2n}C_{2n} + C_{2n-2} = 0$$

$$D_{2n} = \frac{\left[a_x - (2n + \beta_x)^2\right]}{q_x}$$
(3)





Using the eq.(3) recursively one obtains:

$$C_{2n+2} = \frac{C_{2n}}{D_{2n} - \frac{1}{D_{2n+2} - \frac{1}{\dots}}}$$

$$C_{2n} = \frac{C_{2n-2}}{D_{2n} - \frac{1}{D_{2n-2} - \frac{1}{\dots}}}$$

$$\beta_x^2 = a_x - q_x \left( \frac{1}{D_0 - \frac{1}{D_2 - \frac{1}{\dots}}} + \frac{1}{D_0 - \frac{1}{D_{-2} - \frac{1}{\dots}}} \right)$$

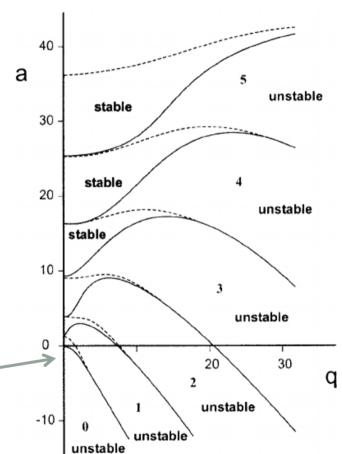
Problem 4: Show that these relations indeed hold





For stability of ion motion  $(x(t \to \infty) < \rho_0)$ 

$$\beta_{x,y,z} = 0/1$$



Lowest order stability region

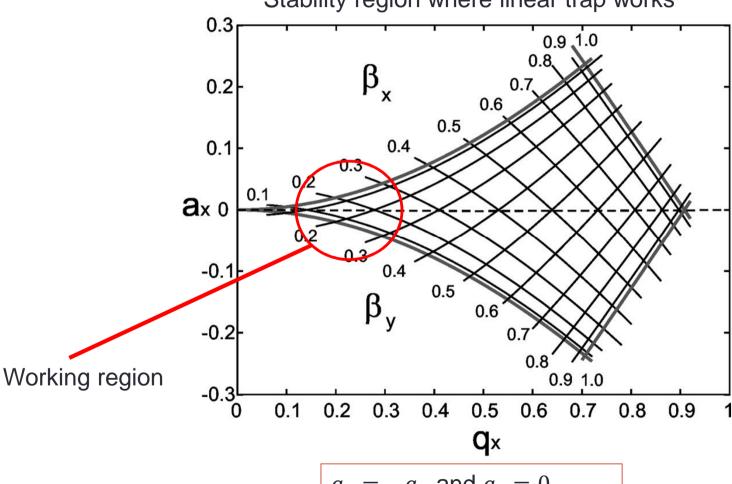
Strutt plot / solutions for 1D Mathieu equation





Unfortunately life is never that simple, "there are instabilities within the stability region"

Stability region where linear trap works



$$q_x = -q_y$$
 and  $q_z = 0$ 





Approximation as per the real linear trap operation:

- Axial confinement is weaker than radial  $a_x < q_x$
- Works close to the origin of 1<sup>st</sup> stability region  $|a_x|$ ,  $q_x^2 \ll 1$
- Assuming  $C_{+4} \approx 0$

One obtains: 
$$\beta_x \approx \sqrt{a_x + \frac{q_x^2}{2}}$$

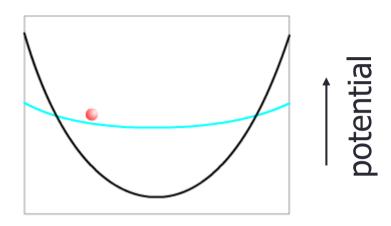
And: 
$$x(t) \approx 2AC_0 \cos\left(\frac{\beta_x \Omega}{2}t\right) \left[1 - \frac{q_x}{2}\cos(\Omega t)\right]$$
  
=  $2AC_0 \cos\left(\frac{\beta_x \Omega}{2}t\right) - \frac{2AC_0 q_x}{2}\cos\frac{\beta_x \Omega}{2}t\cos\Omega t$ 

Secular motion

Micro motion







$$x(t) = 2AC_0 \cos\left(\frac{\beta_x \Omega}{2}t\right) - \frac{2AC_0 q_x}{2} \cos\frac{\beta_x \Omega}{2}t \cos\Omega t$$

Micro motion

Secular motion





### Pseudo-potential approximation

The mean displacement of the ion is negligible within time  $\frac{1}{0}$ 

The total displacement is composed of secular and micro-motion parts

$$x = x_{s} + x_{\mu}$$

reminder 
$$x(t) \approx 2AC_0 \cos\left(\frac{\beta_x\Omega}{2}t\right) \left[1 - \frac{q_x}{2}\cos(\Omega t)\right]$$

Secular displacement is large but frequency is slow as compared to micro-motion

$$x_s \gg x_\mu$$

$$\dot{x}_{s} \ll \dot{x}_{\mu}$$

The time-dependent motion in x is

$$\ddot{x} + (a_x - 2q_x \cos 2\zeta)x = 0$$

$$\ddot{x_{\mu}} = -(a_x - 2q_x \cos 2\zeta) x_s \tag{1}$$

$$a_x \ll q_x$$

Integrating over time

$$x_{\mu} = -\frac{q_x}{2}\cos 2\zeta \ x_s$$

 $x_s$  is constant in one period





### Pseudo-potential approximation

Therefore the amplitude of motion is:

$$x = x_s - \frac{q_x}{2}\cos 2\zeta \ x_s$$

Substituting in  $\ddot{x} + (a_x - 2q_x cos 2\zeta)x = 0$ :

$$\ddot{x} = -(a_x - 2q_x \cos 2\zeta) \left(1 - \frac{q_x}{2} \cos 2\zeta\right) x_s$$

$$= -a_x x_s - q_x^2 \cos^2 2\zeta x_s + 2q_x \cos 2\zeta x_s + \frac{q_x a_x}{2} \cos 2\zeta x_s$$

Averaging over one cycle of RF:

$$<\ddot{x_s}> = -\left(a_x + \frac{q_x^2}{2}\right)x_s$$

$$\left\langle \frac{d^2 x_s}{dt^2} \right\rangle = -\left(a_x + \frac{q_x^2}{2}\right) \frac{\Omega^2}{4} x_s$$

Reminder: 
$$\zeta = \frac{\Omega t}{2}$$





### Pseudo-potential approximation

From pseudo potential model one obtains:

$$\left\langle \frac{d^2 x_s}{dt^2} \right\rangle = -\left(a_x + \frac{q_x^2}{2}\right) \frac{\Omega^2}{4} x_s = -\frac{\beta_x^2 \Omega^2}{4} x_s = -\omega_x^2 x_s$$

From solving the Mathieu equation one obtains:

$$x(t) = 2AC_0 \cos\left(\frac{\beta_x \Omega}{2}t\right) - \frac{2AC_0 q_x}{2} \cos\frac{\beta_x \Omega}{2}t \cos\Omega t$$

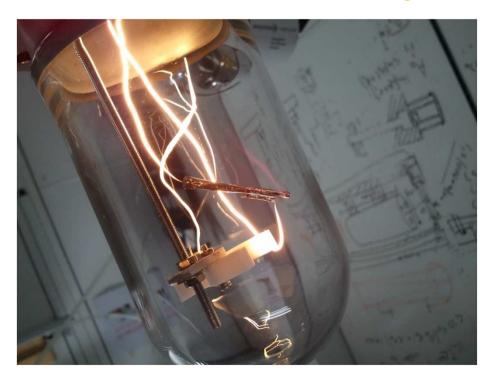
Therefore they match and we observe that the motion is a <u>simple</u> <u>harmonic oscillator motion</u>

Problem 5: Prove that the pseudo-potential trap depth is  $\overline{D}_{\chi} = \frac{eV_0^2}{4mr_0^2\Omega^2}$  considering  $a_{\chi} = 0$ .





#### Atomic oven – resistive heating

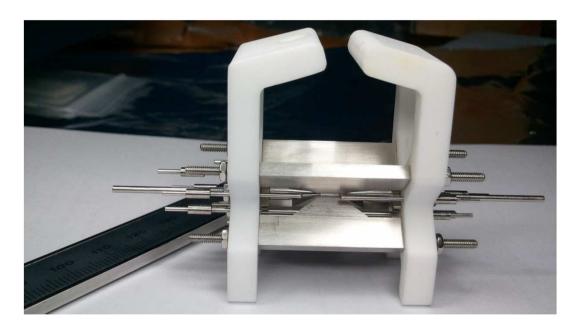








#### Trap assembly – UHV protocols

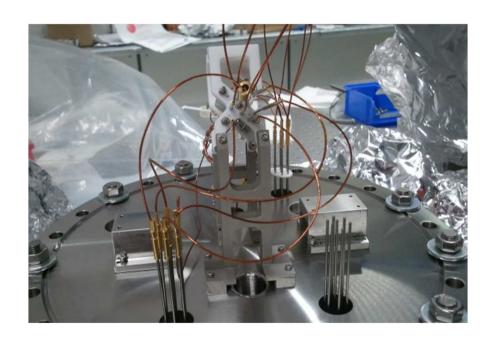


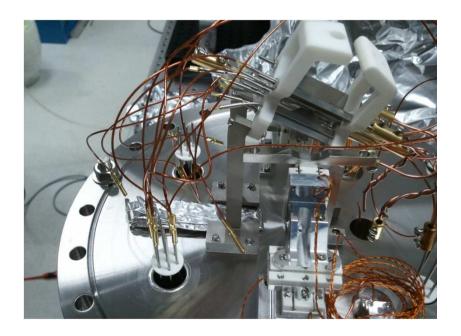






Trap assembly – UHV protocols





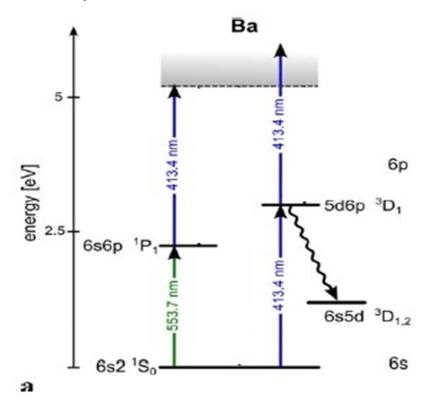




#### Ion creation - in-situ

- 1. Electron impact
- 2. Surface ionization
- 3. Resonant laser ionization
- 4. etc.

#### Example for Ba+







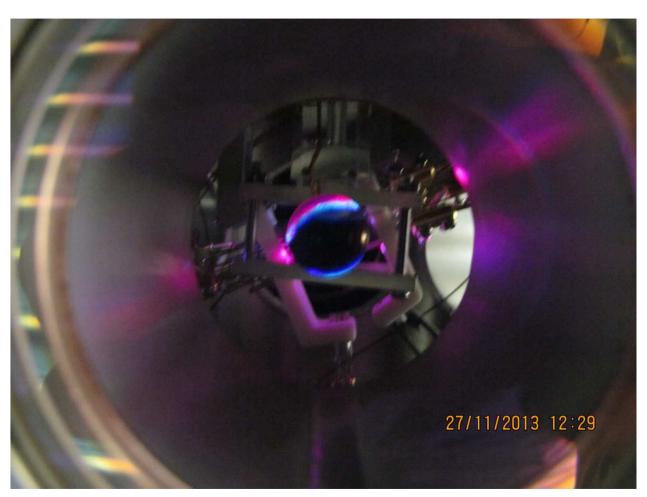
### Ion trap drive





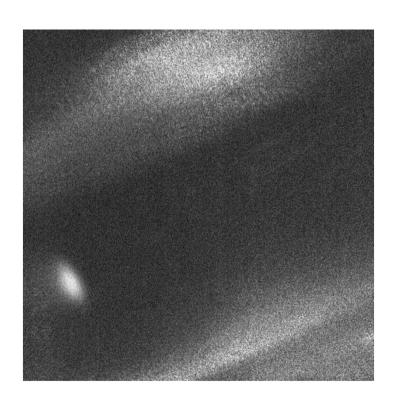


Ion imaging





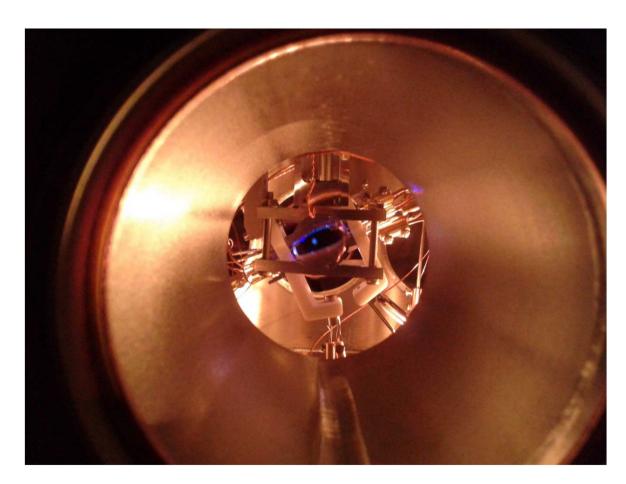




Is there ion?



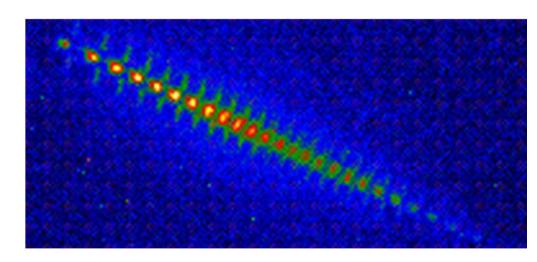




Smart phone image – not Apple!!







Chain of 25 ions CCD image